UNIVERSITY OF SWAZILAND 136 **FACULTY OF SCIENCE DEPARTMENT OF PHYSICS** MAIN EXAMINATION 2011/2012 TITLE OF PAPER **ELECTROMAGNETIC THEORY** I : **COURSE NUMBER** P331 : **TIME ALLOWED THREE HOURS** : **INSTRUCTIONS** ANSWER ANY FOUR OUT OF FIVE : **QUESTIONS. EACH QUESTION CARRIES 25** MARKS. MARKS FOR DIFFERENT SECTIONS

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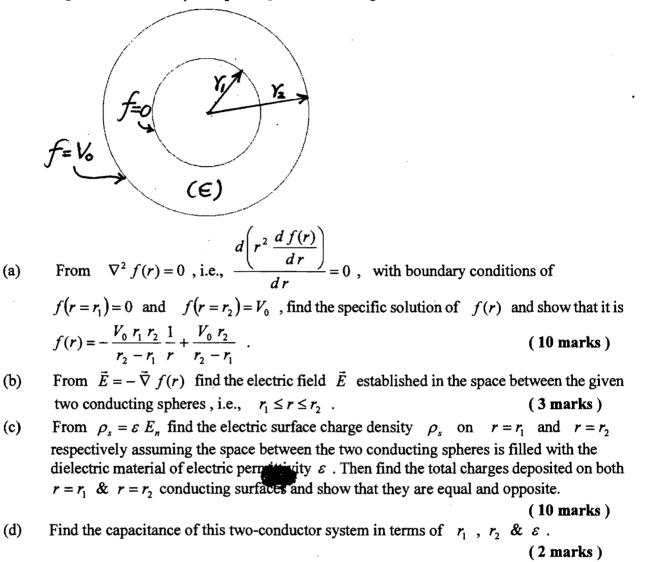
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P331 Electromagnetic Theory I

Question one

A potential difference V_0 is maintained between two very thin-shelled and co-centred spherical conducting balls of radius $r_1 \& r_2$ as depicted in the diagram below :



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Question two

- (a) A very long thin conducting wire is situated at z-axis and uniformly charged with line charge density σ_i .
 - (i) Use integral Coulomb's Law and choose an appropriate Gaussian surface to deduce that the electric field at a field point outside the given thin conducting wire is

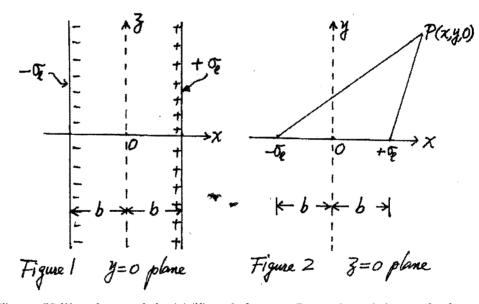
$$\vec{E} = \vec{e}_{\rho} \frac{\sigma_l}{2 \pi \varepsilon_0 \rho}$$
 where ρ is the distance from z-axis and

 \vec{e}_{o} is one of the unit vectors in cylindrical coordinate system (7 marks)

(ii) Use $\Phi = -\int_{P_0}^{P} \vec{E} \cdot d\vec{l}$ to find the electric potential at any point $P:(\rho,\phi,z)$ where $P_0:(\rho_0,0,0)$ is the zero potential reference point and show that it is

$$\Phi = \frac{\sigma_l}{2 \pi \varepsilon_0} \ln \left(\frac{\rho_0}{\rho} \right)$$
 (5 marks)

(b) Two long thin conducting wires are parallel to the z-axis and lie in the y = 0 plane, i.e, in the x-z plane. One is situated at x = -b and carries $-\sigma_i$ uniform line charge density and the other is situated at x = +b and carries $+\sigma_i$ uniform line charge density, as shown in the Figure 1 (on y = 0 plane) and Figure 2 (on z = 0 plane) below :



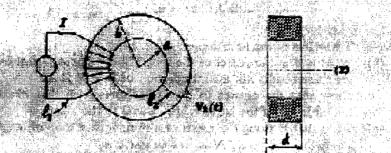
(i) Utilize the result in (a)(ii) and choose P_0 as the origin, apply the superposition principle to deduce that the electric potential at point P:(x, y, 0) is

$$\Phi = \frac{\sigma_l}{4 \pi \varepsilon_0} \ln \left(\frac{(x+b)^2 + y^2}{(x-b)^2 + y^2} \right)$$
 (7 marks)

(ii) Use $\vec{E} = -\vec{\nabla} \Phi$ to find the electric field \vec{E} generated by the given two conducting wires. Also find the value of \vec{E} at the origin. (4 + 2 marks)

Question three

(a) A static current I flows in the n_1 turn toroid l_1 wired around an iron ring core of magnetic permeability μ with the rectangular cross-section area $(b-a) \times d$ as shown below



- (i) Use the integral Ampere's law and choose proper closed loops to find \vec{B} in terms of ρ , n_1 , μ & I within the iron core, i.e., $a \le \rho \le b$ & $0 \le z \le d$ region. (7 marks)
- (ii) Find the total magnetic flux Ψ_m passing through the cross-section area $(b-a) \times d$ of the iron ring in counter clockwise sense, i.e., $\int_{S} \vec{B} \cdot d\vec{s}$ where $S: a \le \rho \le b$, $0 \le z \le d$ & $d\vec{s} = \vec{a}_{\phi} d\rho dz$, in terms of a, b, d, n_1, μ & I. (6 marks)
- (iii) Find the external self-inductance L_e of the toroid wire l_1 in terms of $a, b, d, \mu \& n_1$. (3 marks)
- (b) Placing a single turn secondary coil l_2 around the iron ring and if the toroid wire l_1 carries a sinusoidal current of $I_0 \sin(\omega t)$ instead of carrying a static current I, find the induced e.m.f. $V_2(t)$ for the single turn secondary coil l_2 in terms of $a, b, d, \omega, n_1, \mu \& I_0$ under quasi static situation. If $a = 4 \ cm$, $b = 8 \ cm$, $d = 2 \ cm$, $n_1 = 100$, $f = 50 \ Hz$, $\mu = 400 \ \mu_0$ and $I_0 = 2 \ A$, compute the amplitude of $V_2(t)$. (6+3 marks)

Question four

(a) (i) From the time-dependent Maxwell's equations deduce the following wave equation for \vec{E} in the material region with parameters of μ , $\varepsilon \& \sigma$ where $\rho_v = 0 \& \vec{J} = \sigma \vec{E}$, as

$$\nabla^2 \vec{E} = \mu \sigma \frac{\partial \vec{E}}{\partial t} + \mu \varepsilon \frac{\partial^2 \vec{E}}{\partial t^2}$$
 (6 marks)

(ii) By direct substitution, show that $E_x = \hat{E}_m e^{\hat{\gamma} \cdot z} e^{i\omega t}$ (where \hat{E}_m is any constant, ω is any frequency and $\hat{\gamma} = \sqrt{i\omega\mu\sigma - \omega^2\mu\varepsilon}$) is a solution to the E_x part of the wave equation in (a)(i), i.e., $\nabla^2 E_x = \mu \sigma \frac{\partial E_x}{\partial t} + \mu \varepsilon \frac{\partial^2 E_x}{\partial t^2}$. (6 marks)

- (b) An uniform plane wave traveling along the +z direction with the field components $E_x(z) \& H_y(z)$ has a complex electric field amplitude $\hat{E}_m = 200 e^{i 40^0} \frac{V}{m}$ and propagates at a frequency $f = 10^7$ Hz in a material region having the parameters $\mu = \mu_0$, $\varepsilon = 4 \varepsilon_0 \& \frac{\sigma}{\omega \varepsilon} = 1$.
 - (i) Find the values of the propagation constant $\hat{\gamma} (= \alpha + i \beta)$ and the intrinsic wave impedance $\hat{\eta}$ for this wave, (5 marks)
 - (ii) Express the electric and magnetic fields in both their complex and real-time forms, with the numerical values of (b)(i) inserted, (5 marks)
 - (iii) Find the values of the penetration depth, wavelength and phase velocity of the given wave . (3 marks)

Question five

An uniform plane wave $(\hat{E}_{x1}^+, \hat{H}_{y1}^+)$, operating at a frequency f, is normally incident upon a layer of d_2 thickness, and emerges to region 3 as shown below :

 0_1 , 0_2 & 0_3 are the respective origins for region 1, 2 & 3 chosen at the first and second interface.

(a) Define for the i^{th} region (i = 1, 2, 3) the reflection coefficient $\hat{\Gamma}_i(z)$ and the total wave impedance $\hat{Z}_i(z)$ and deduce the following :

$$\begin{cases} \hat{Z}_{i}(z) = \hat{\eta}_{i} \frac{1 + \hat{\Gamma}_{i}(z)}{1 - \hat{\Gamma}_{i}(z)} \\ \hat{\Gamma}_{i}(z') = \hat{\Gamma}_{i}(z) e^{2\hat{\gamma}_{i}(z'-z)} & \text{where } z' \& z \text{ are two positions in } i'' \text{ region} \end{cases}$$
(5 + 4 marks)

(b) If $f = 10^5$ Hz & $d_2 = \frac{\lambda_2}{4}$, region 1 & 3 are air regions and region 2 is a lossless

region with parameters $\mu_2 = \mu_0$, $\varepsilon_1 = 9 \varepsilon_0 \& \frac{\sigma}{\omega \varepsilon} = 0$

(i) find the values of β_1 , β_2 , β_3 , λ_2 & $\hat{\eta}_2$, (note: $\hat{\eta}_1 = \hat{\eta}_3 = 120 \pi \Omega$ and $\alpha_1 = \alpha_2 = \alpha_3 = 0$) (4 marks)

(ii) starting with $\hat{\Gamma}_3(z) = 0$ for the rightmost region, i.e., region 3, and using continuous \hat{Z} at the interface as well as the equations in (a), find the values of $\hat{Z}_3(0)$, $\hat{Z}_2(0)$, $\hat{\Gamma}_2(0)$, $\hat{\Gamma}_2(-10 \ cm)$, $\hat{Z}_2(-10 \ cm)$, $\hat{Z}_1(0)$ & $\hat{\Gamma}_1(0)$ (10 marks)

(iii) find the value of
$$\hat{E}_{m1}^-$$
 if given $\hat{E}_{m1}^+ = 100 e^{i 50^0} \frac{V}{m}$ (2 marks)

Useful informations

$$e = 1.6 \times 10^{-19} C$$

$$m_e = 9.1 \times 10^{-31} kg$$

$$\mu_0 = 4 \pi \times 10^{-7} \frac{H}{m}$$

$$\varepsilon_0 = 8.85 \times 10^{-12} \frac{F}{m}$$

$$\hat{r} = \alpha + i\beta \qquad \text{where}$$

$$\alpha = \frac{\omega \sqrt{\mu \varepsilon}}{\sqrt{2}} \sqrt{\sqrt{1 + \left(\frac{\sigma}{\omega \varepsilon}\right)^2} - 1}$$

$$\beta = \frac{\omega \sqrt{\mu \varepsilon}}{\sqrt{2}} \sqrt{\sqrt{1 + \left(\frac{\sigma}{\omega \varepsilon}\right)^2} + 1}$$

$$\frac{1}{\sqrt{\mu_0 \varepsilon_0}} = 3 \times 10^8 \frac{m}{s}$$

$$\hat{\eta} = \frac{\sqrt{\frac{\mu}{\varepsilon}}}{\sqrt{1 + \left(\frac{\sigma}{\omega \varepsilon}\right)^2}} e^{i\frac{1}{2}\tan^{-1}\left(\frac{\sigma}{\omega \varepsilon}\right)}$$

$$\eta_0 = 120 \pi \ \Omega = 377 \ \Omega$$

$$\beta_0 = \omega \sqrt{\mu_0 \varepsilon_0}$$

$$\oiint_s \vec{E} \cdot d\vec{s} = \frac{1}{\varepsilon} \iiint_{\nu} \rho_{\nu} d\nu$$

$$\oiint_s \vec{B} \cdot d\vec{s} = 0$$

$$\oint_L \vec{E} \cdot d\vec{l} = -\frac{\partial}{\partial t} \left(\iint_s \vec{B} \cdot d\vec{s}\right)$$

$$\begin{split} \vec{\nabla} \bullet \vec{E} &= \frac{\rho_{-}}{e} \\ \vec{\nabla} \bullet \vec{B} &= 0 \\ \vec{\nabla} \times \vec{E} &= -\frac{\partial \vec{B}}{\partial t} \\ \vec{\nabla} \times \vec{B} &= \mu \vec{J} + \mu \varepsilon \frac{\partial \vec{E}}{\partial t} \\ \vec{J} &= \sigma \vec{E} \\ \underbrace{ \iint_{S} \vec{F} \bullet d\vec{s} &= \underbrace{ \iint_{K} (\vec{\nabla} \cdot \vec{F}) d\nu }_{S} \quad divergence theorem \\ \underbrace{ \iint_{S} \vec{F} \bullet d\vec{s} &= \underbrace{ \iint_{K} (\vec{\nabla} \cdot \vec{F}) \bullet d\vec{s} }_{S} \quad Stokes' theorem \\ \vec{\nabla} \cdot (\vec{\nabla} \times \vec{F}) &= 0 \\ \vec{\nabla} \times (\vec{\nabla} \times \vec{F}) &= \vec{\nabla} (\vec{\nabla} \bullet \vec{F}) - \nabla^{2} \vec{F} \\ \vec{\nabla} f &= \vec{a}_{*} \frac{\partial f}{\partial t} + \vec{e}_{*} \frac{\partial f}{\partial t} + \vec{e}_{*} \frac{\partial f}{\partial z} \\ &= \vec{e}_{*} \frac{\partial f}{\partial t} + \vec{e}_{*} \frac{1}{2} \frac{\partial f}{\partial y} + \vec{e}_{*} \frac{\partial f}{\partial z} \\ = \vec{e}_{*} \frac{\partial f}{\partial t} + \vec{e}_{*} \frac{1}{2} \frac{\partial f}{\partial y} + \vec{e}_{*} \frac{\partial (F_{*})}{\partial z} \\ &= \frac{1}{\rho} \frac{\partial (F_{*})}{\partial x} + \frac{\partial (F_{*})}{\partial z} \\ &= \frac{1}{\rho} \frac{\partial (F_{*})}{\partial x} + \frac{\partial (F_{*})}{\partial z} \\ &= \frac{1}{\rho} \frac{\partial (F_{*})}{\partial x} + \frac{\partial (F_{*})}{\partial z} \\ &= \frac{1}{\rho} \frac{\partial (F_{*})}{\partial y} - \frac{\partial (F_{*})}{\partial z} \\ &= \frac{1}{\rho} \frac{\partial (F_{*})}{\partial \phi} - \frac{\partial (F_{*})}{\partial z} \\ &= \frac{1}{\rho} \frac{\partial (F_{*} \sin(\theta))}{\partial \phi} \\ &= \frac{\vec{e}_{*}}{\rho} \left(\frac{\partial (F_{*} r \sin(\theta))}{\partial \phi} \\ &= \frac{\partial (F_{*} r)}{\rho} \\ &= \frac{\vec{e}_{*}}{\rho} \left(\frac{\partial (F_{*} r \sin(\theta))}{\partial \theta} \\ &= \frac{\partial (F_{*} r)}{\partial z} \\ &= \frac{\partial (F_{*} r)}{\rho} \\ &= \frac{\partial (F_{*} r)}{\partial \phi} \\ &= \frac{\partial (F_{*} r)}{\rho} \\ \\ &= \frac{\partial (F_{*} r)}{\rho} \\ \\ &= \frac{\partial (F_{*} r)}{\rho} \\ \\ &= \frac{\partial$$

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